

# Black Hole Entropy in String Theory ...

Donald Marolf  
Syracuse University

APS meeting,  
April 07, 2002

...and answers to related questions.

# Summary

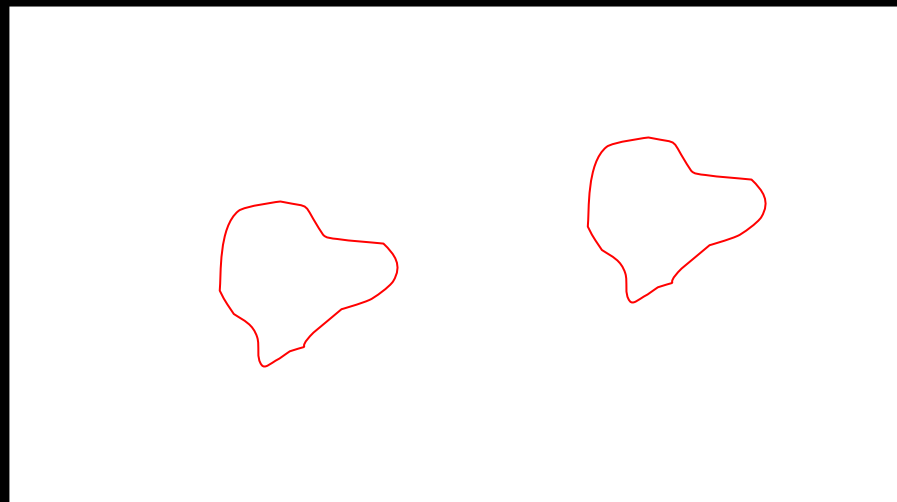
- String Theory: Under construction\*, but contains classical supergravity (and black holes) in limit of low curvatures.
  - Contains black holes, and can calculate entropy for *some*. Find:  $S = A/4G$
  - However, no general results explaining why  $S \sim A$ .
  - Other conclusive results to come?\*
- \* = Juan may comment from ADS/CFT.

# The setting, and a question:

Black hole is a nonperturbative curved spacetime description.

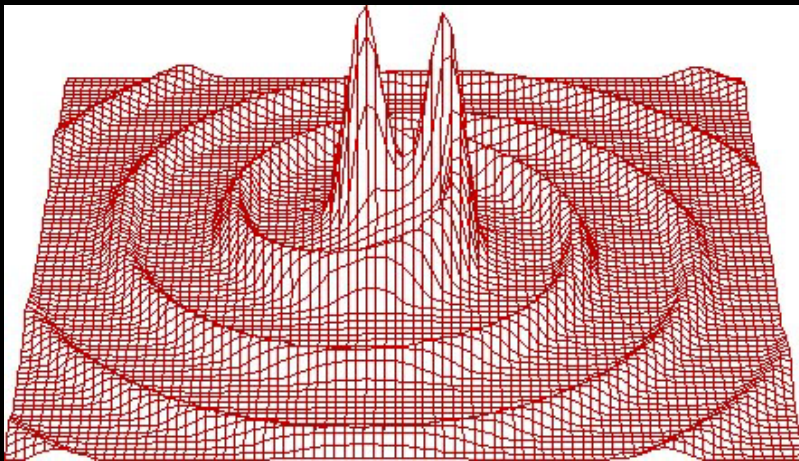
vs.

Most powerful/most developed tool in string theory:  
perturbation theory about flat spacetime –  
i.e., strings in flat space.



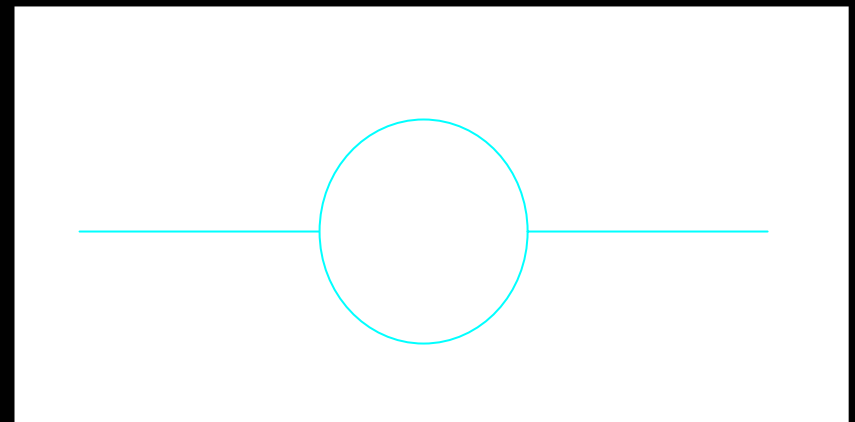
# Answer (v1):

- Finite Curvature does lead to a tension, but need only require hard work.
- Very Old Story (Feynmann, DeWitt, etc):



(LSC)

=



Gravitons (& more)  
on flat background...

And gravitons ARE strings!

But BH's are non-perturbative?

Answer (v2): ``Luck!''

Supergravity Warning: The Following Explanation  
has been altered to fit the audience.  
Many U(1)  $\longrightarrow$  Analogues of R-N  
Maxwell fields

'Backwards' = anti-historical  
Similar Hawking Emission

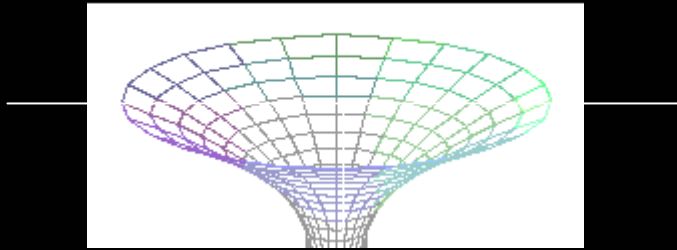
But, SUSY  $\longrightarrow$  BPS  $\longrightarrow$   $Q < M$

$(M-Q)_{\text{emitted}} > 0$ , radiates toward extremality if  $Q \sim M$ .

Can make extreme black holes!

# Connection to perturbation theory:

Take  $MG \rightarrow 0$



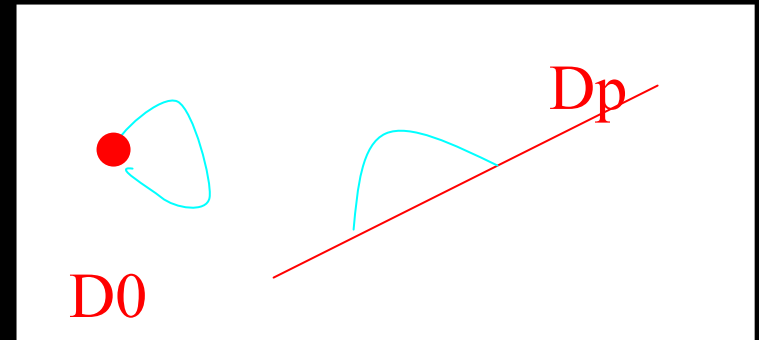
Nearly flat, consistent with perturbative construction if can find new 'ingredient' to represent object at origin.

Note: Large Curvatures at 'origin', Quantum effects important

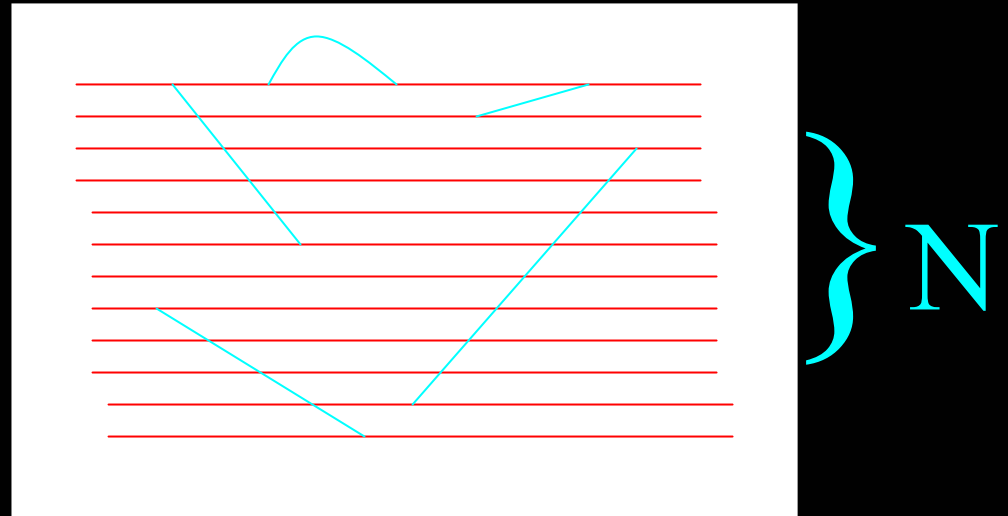
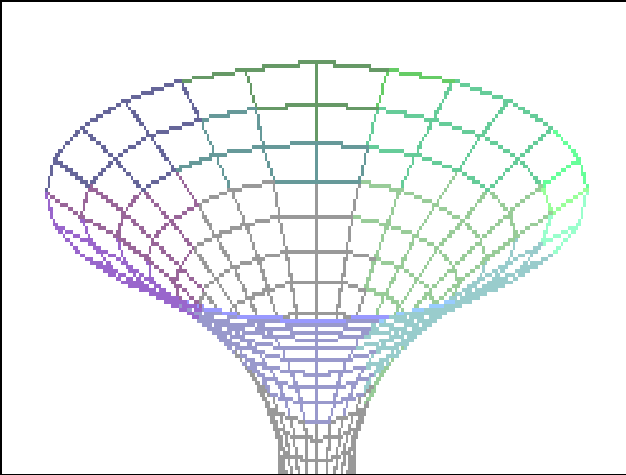
Guess:

D-branes, D0-brane

w/ open strings =  
excitations of black brane



# Or bound state of N branes for larger BH:

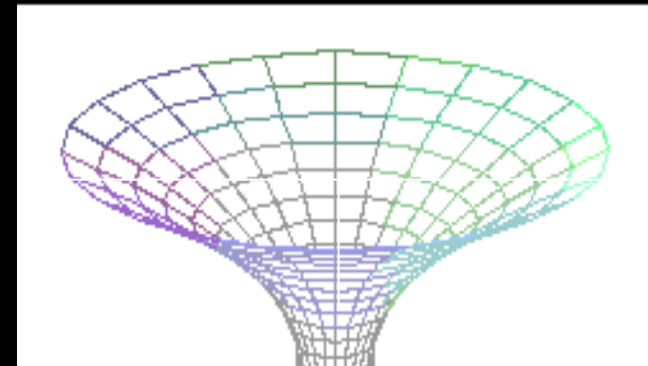
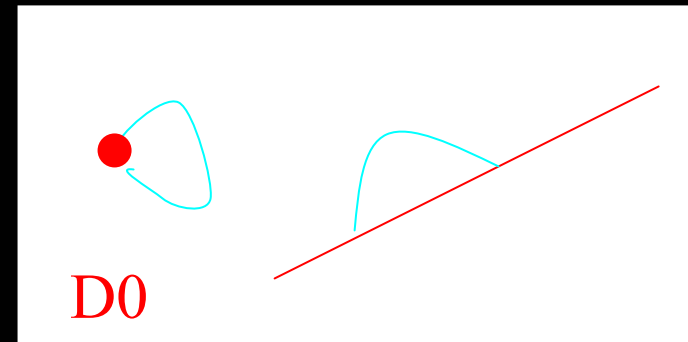


Note:  $O(N^2)$  degrees of freedom.

# Many checks:

- Same charges, symmetries (SUSY)
- Near extremality open strings give correct dynamics for bound states, including coupling to other modes.
- Why such agreement? SUSY\*.
- Far from extremality = hard to calculate in BH regime. Strongly coupled QFT\*.

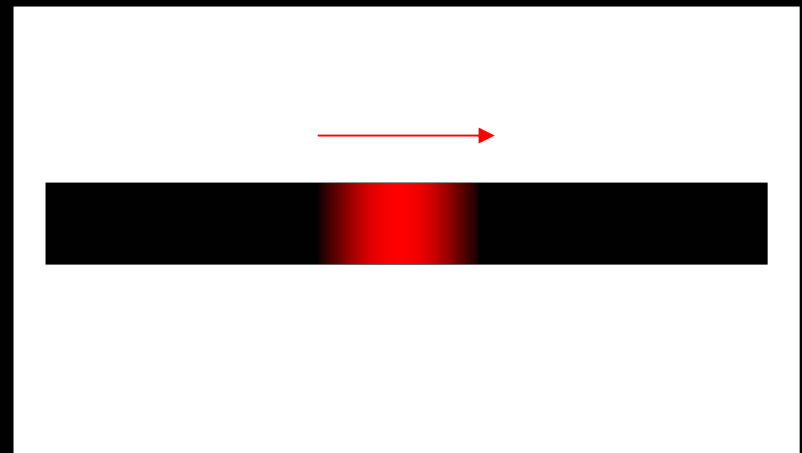
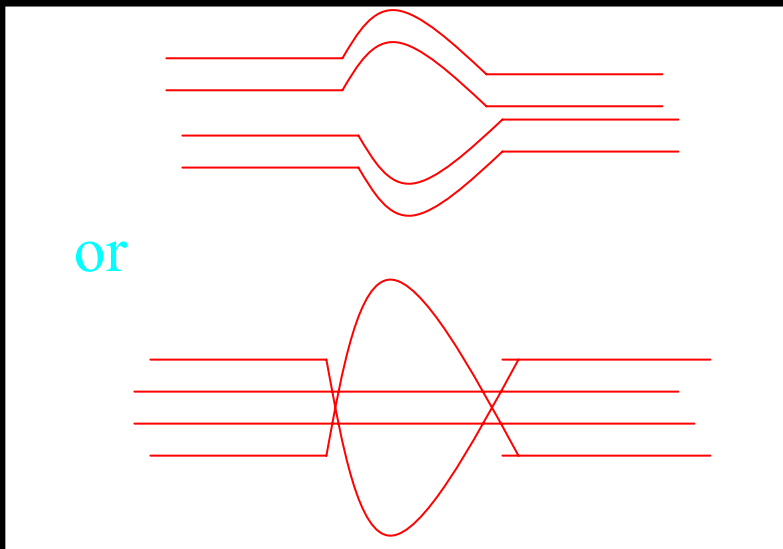
\* = More from Juan?





# But one difference!

- An excitation of the BH maps to a *class* of D-brane states! Degeneracy = entropy!
- Large (classical) BH  $\longrightarrow$  bound state of many D-branes. Black hole seems to be 'coarse graining' over D-brane states.
- E.g. total momentum density – not distribution.



# Classic Results:

Strominger & vafa + successors

Getting large smooth horizon for extreme BH is tricky  $\rightarrow$  use 3 (or 4) types of charge:  $Q_e$ ,  $Q_m$ , and  $P$ .

$$10d: M = R^5 \times S^1 \times T^4$$

$$8d \text{ horizon: } M = S^3 \times S^1 \times T^4$$

$$Q_e Q_m = \text{integer (Dirac)}$$

$$N = PL/2\pi\hbar$$

$$S = 2\pi \sqrt{Q_e Q_m N}$$

D-brane description has degenerate ground state with same  $S$  to leading order.

# (Some) Further results

Entropy for other SUSY  
and near SUSY BH's and  
branes

Balasubramanian,  
Breckenridge, Callan,  
Cvetic, Horowitz,  
Johnson, Larsen, Lowe,  
Khuri, Klebanov,  
Maldacena, Marolf,  
Meyers, Peet,  
Strominger, Tseytlin,  
+ ...

Emission/Absorption for  
long  $\lambda$  near SUSY BH's  
and branes

Das & Mathur  
Dhar, Mandal, & Wadia  
Maldacena & Strominger  
(grey body factors) +...

# Other Results

<p>Corrections to <math>S=A/4</math></p> $S = \frac{1}{g} (R + R^2 + \dots)$ <p>for SUSY BH's</p>	<p>G. L. Cardoso, T. Mohaupt, &amp; B. de Wit</p>
<p>Non-extremal BTZ</p>	<p>Complicated history.</p>
<p>A certain non-extreme black hole in a funny 10-dim space.</p> $\beta F \sim -4.1 T^{1.8}$	<p>G. Lifschytz, D. Lowe, D. Kabat – numerical mean field theory for QFT. Also uses ADS/CFT-like duality.</p> $\beta F \sim -2 T^{1.7}$

# Summary

1. String theory provides a remarkably consistent story which correctly predicts BH entropy in many situations – finds famous factors of  $\frac{1}{4}$  in  $S = A/4$ ..
2. But (so far) yields little insight into form of  $S$ , or spacetime description ... Why  $A$ ? Why always  $\frac{1}{4}$ ?
3. Computations are prohibitively hard for many interesting cases (e.g., Schwarzschild).